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Nonlinear Viscoelastic Behaviour of *Brain Tissue*

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Abstract

In this worksheet the relaxation of *Brain Tissue* has been calculated to experimental data by two nonlinear model functions:

a **five parameters** *PRONY*-Series and compared with a **three parameters** *Sqrt(t)*-Law due to *BETTEN, Creep Mechanics*, 3rd edition, 2008 Springer-Verlag Berlin / Heidelberg.

The experiments have been carried out by *P. AMEDIEU* and published in his dissertation *Contribution à la biomécanique des tissus mous intracrâniens*, 2004 Université de Picardie Jules Verne, Faculte de Medicine, Amiens Cedex France

Relaxation

```
[ > restart:  
[ > with(Statistics):  
[ > X:=vector([0,25,50,100,150,200,250]);  
[ X := [0, 25, 50, 100, 150, 200, 250]  
[ > whattype(X);  
[ symbol  
[ > Y:=vector([1,0.57,0.49,0.426,0.398,0.383,0.374]);  
[ Y := [1, 0.57, 0.49, 0.426, 0.398, 0.383, 0.374]  
[ > whattype(Y);  
[ symbol  
[ > DATA:=seq([X[i],Y[i]],i=1..7);  
[ DATA :=  
[ [0, 1], [25, 0.57], [50, 0.49], [100, 0.426], [150, 0.398], [200, 0.383], [250, 0.374]
```

PRONY-Series:

```
[ > R(t):=evalf(NonlinearFit(A+B*exp(-C*t)+D*exp(-E*t),X,Y,t),4);  
[ R(t) := 0.3674 + 0.3937 e(-0.09820 t) + 0.2389 e(-0.01388 t)  
[ > R(0):=evalf(subs(t=0,R(t)));
```

```
R(infinity):=simplify(subs(t=infinity,R(t)));
```

```
R(0) := 1.0000
```

```
R(∞) := 0.3674
```

Sqrt(t)-Law:

```
> r(t):=evalf(NonlinearFit(a+b*exp(-c*sqrt(t)),X,Y,t),4);
```

$$r(t) := 0.3540 + 0.6461 e^{(-0.2196\sqrt{t})}$$

```
> r(0):=simplify(subs(t=0,r(t)));
```

```
r(infinity):=simplify(subs(t=infinity,r(t)));
```

```
r(0) := 1.0001
```

```
r(∞) := 0.3540
```

```
> alias(th=thickness,co=color):
```

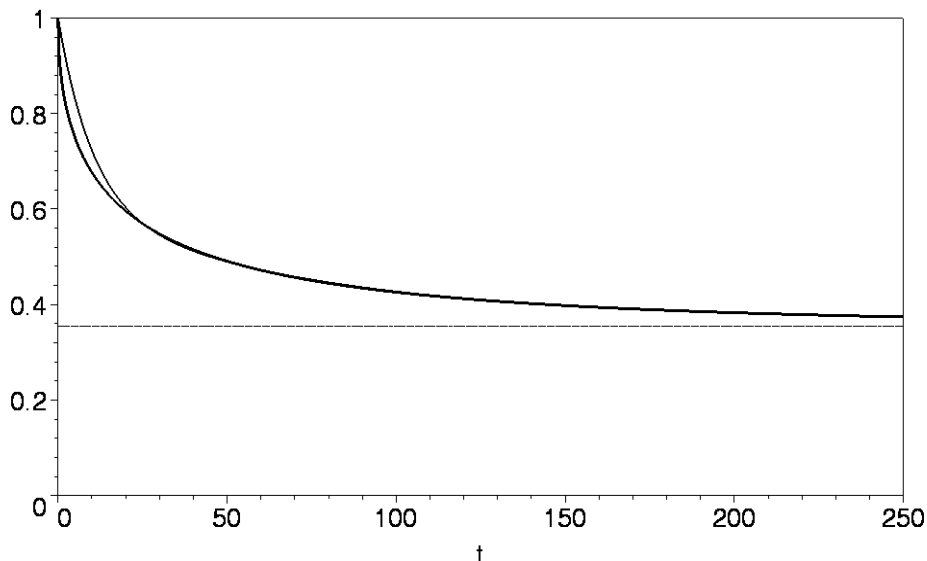
```
> p[1]:=plot(R(t),t=0..250,co=black,th=2,axes=boxed):
```

```
> p[2]:=plot(r(t),t=0..250,0..1,co=black,th=3,  
title="Five-Parameter PRONY & Three-Parameter Sqrt(t)");
```

```
> p[3]:=plot(0.354,t=0..250,co=black,linestyle=4):
```

```
> plots[display](seq(p[k],k=1..3));
```

Five-Parameter PRONY & Three-Parameter Sqrt(t)



```
> # L[2]-Distance-Norm between PRONY R(t) and Sqrt(t)-Law r(t):
```

```
> L[2]:=sqrt((1/250)*Int((R(tau)-r(tau))^2,tau=0..250))=  
evalf(sqrt((1/250)*int((R(t)-r(t))^2,t=0..250)),4);
```

$$L_2 := \frac{1}{50} \sqrt{10} \sqrt{\int_0^{250} (R(\tau) - r(\tau))^2 d\tau} = 0.01623$$

Error-Norm for PRONY R(t) to Experimental DATA:

```
> with(linalg):
```

```
> for i from 1 to 7 do
```

```
  v[i]:=evalf(subs(t=DATA[i][1],R(t))-DATA[i][2]) od:
```

```
> V:=vector([seq(v[i],i=1..7)]);
V := [0., 0.0000601329, -0.0003492771, 0.00104461655, -0.00081364654, -0.00071959670,
      0.000833858634]
```

```
> L[2]:=(1/sqrt(number_of_points))*Norm(V,2)=
evalf((1/sqrt(7))*norm(V,2),4);
```

$$L_2 := \frac{\text{Norm}(V, 2)}{\sqrt{\text{number_of_points}}} = 0.0006649$$

Error-Norm for $\text{Sqrt}(t)$ -Law $r(t)$ to Experimental DATA:

```
> for i from 1 to 7 do
w[i]:=evalf(subs(t=DATA[i][1],r(t))-DATA[i][2]) od:
```

```
> W:=vector([seq(w[i],i=1..7)]);
```

```
W := [0.0001, -0.0005014264, 0.0007490916, -0.00012314622, -0.00012193427,
      -0.00005662576, 0.00006091550]
```

```
> l[2]:=(1/sqrt(number_of_points))*Norm(W,2)=
evalf((1/sqrt(7))*norm(W,2),4);
```

$$l_2 := \frac{\text{Norm}(W, 2)}{\sqrt{\text{number_of_points}}} = 0.0003503$$

```
> Q:=evalf(L[2]/l[2],4);
```

$$Q := \frac{\text{Norm}(V, 2)}{\text{Norm}(W, 2)} = 1.898$$

```
>
```

The above error-norms $L[2]$, $l[2]$, or $Q = L[2] / l[2]$ illustrate that the $\text{Sqrt}(t)$ -Law with only three parameters is a better approximation as the PRONY-Series with five parameters to be determined. In this worksheet and in a lot of other examples of *stress relaxation* or even in *creep problems* BETTEN has analysed in more detail that the $\text{SQRT}(t)$ -Law furnishes the best approximation.

Stress and Structural Relaxation

Besides the *stress relaxation* another kind, namely the *structural relaxation*, is also important. This sort of relaxation governs the time-dependent response of a fluid to a change in temperature. For instance, a liquid is held at temperature $T[1]$ until the property $p(t, T[1])$ reaches its equilibrium value $p(\text{infinity}, T[1])$, then it is suddenly cooled to $T[2]$. The instantaneous change in p is proportional to the difference $T[1] - T[2]$, followed by relaxation toward the equilibrium value $p(\text{infinity}, T[2])$.

Based upon a lot of experiments on *glass*, SCHERER, G. (1986), *Relaxation in Glass and Composites*, has shown that both the *stress* and *structural relaxation* in glass can be predicted by the relation

```
> restart:
```

```
> r(t):=exp(-(t/lambda)^b);
```

$$r(t) := e^{-\left(\frac{t}{\lambda}\right)^b}$$

often called *KOHLRAUSCH function*, which is a modified form of the *MAXWELL* relaxation function with $b = 1$. The exponent b for a variety of glasses was found by *SCHERER* to be near the value of $b = 0.5$, so that the assumption of the *Sqrt(t)-Law*, introduced by *BETTEN*, is justified. However, the relation $r(t)$ is valid only for *stabilized glass*, i.e., the glass is held at a given temperature until its properties do no longer change with time, then the load can be applied. In unstabilized glasses the viscosity η and other typical properties, e.g., the density, vary with time. Then, the relaxation function should be replaced by the formular

> **restart:**

> **`r(t) := exp(- (Int(G[0]/eta(tau), tau=0..t))^b);`**

$$r(t) := e^{-\left(\int_0^t \frac{G_0}{\eta(\tau)} d\tau\right)^b}$$

where, in agreement with experimental results, the exponent b can again be assumed to $b = 0.5$, as has been discussed by *SCHERER* in more detail.

>